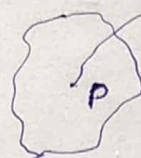


Electric Potential

● Electrostatic Potential :-

Workdone to bring a unit positive charge from infinity to any point in the electric field is called electric potential.



OR. The workdone in moving a unit positive charge from infinity to that point against the electrostatic force of the electric field of charge irrespective of the path followed.

Now, Potential $(V)_r = \frac{W}{q}$; if $q = 1C$, then $V = W$

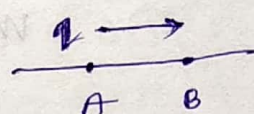
Unit: Potential (V) = $\frac{1J}{1C}$ or, 1V. Unit of Potential = Volt.

Note $V_p - V_\infty = \frac{W}{q}$ [difference betⁿ ∞ and P]
 here, $V_\infty \rightarrow 0$

$V_p = \frac{W}{q}$ \therefore Electric potential is a scalar quantity.

● Potential Difference :-

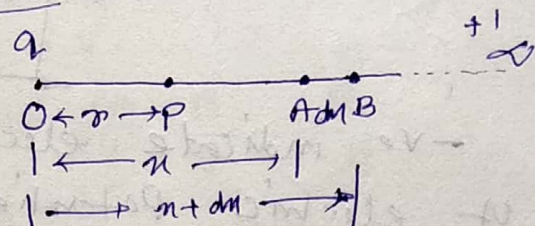
$V_A - V_B = \Delta V = \frac{W_{AB}}{q}$



if $q = 1C$, $\Delta V = W_{AB}$

● Electric Potential due to a point charge :-

Let consider a point charge placed at O. We want to calculate potential at P at a distance r from O. Let us take elementary portion A and B at a distance x and $x+dx$ from O.



Now, electrostatic force of repulsion betⁿ q and $+1$ charge at distance x .

$F = \frac{1}{4\pi\epsilon_0} \cdot \frac{qx}{x^2}$ — (i)

Small workdone to displace charge $+1$ through dx .

$dW = \vec{F} \cdot d\vec{x} = F dx \cos 180^\circ = -F dx$ — (ii)

from eqⁿ (i) & (ii) we get

$dW = -\frac{1}{4\pi\epsilon_0} \cdot \frac{q}{x^2} dx$ — (iii)

now, net work done from ∞ to P.

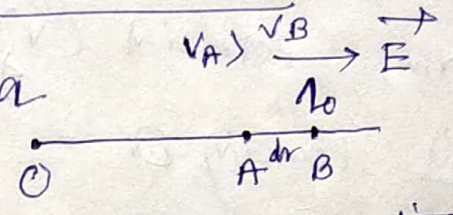
$$\int dw = \int_{x=\infty}^{\infty} - \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r^2} dr$$

$$W = V = - \frac{q}{4\pi\epsilon_0} \left[-\frac{1}{r} \right]_{\infty}^{\infty} = \frac{q}{4\pi\epsilon_0 r}$$

$$\therefore V = \frac{q}{4\pi\epsilon_0 r}$$

Relation betⁿ electric field and potential

Force acting on point B due charge q_0 a

$$\vec{F} = q_0 E \quad \text{--- (i)}$$


now, work done to move charge q_0 from B to A through ~~distance~~

$$dw = \vec{F} \cdot d\vec{r} = F dr \cos 180^\circ = -F dr \quad \text{--- (ii)}$$

from eqⁿ (i) & (ii)

$$dw = -q_0 E dr$$

or, $\frac{dw}{q_0} = -E dr$

or, $dV = -E dr$

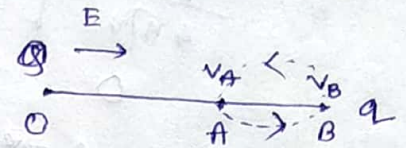
$$\left[dV = \frac{dw}{q_0} \right]$$

$$E = - \frac{dV}{dr}$$

'-ve' indicate electric field intensity is the direction of electric potential decrease.

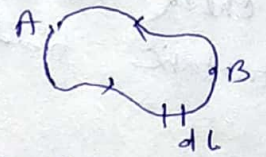
or, electric potential increase to the opposite direction of electric field intensity.

Electrostatic force is conservative:



For conservative force;

(i) Work done by the ~~any~~ force depends on initial and final position ~~does not~~ depends on path.



(ii) Work for the closed ~~surface~~ is zero.

Work done in moving charge q from B to A; $W_{BA} = q_0 (V_A - V_B)$ — (1)

for A to B, $W_{AB} = q_0 (V_B - V_A)$ — (2)

Now from eq^s (1) & (2)

$$W_{BA} + W_{AB} = q_0 [V_A - V_B + V_B - V_A] = 0$$

Hence, for closed line integral $\oint \vec{E} \cdot d\vec{l} = W_{BA} + W_{AB} = 0$ — (3)

According to Stoke's law, $\nabla \times \vec{E} = 0$ — (4)

Hence electrostatic field is conservative.

Poisson's and Laplace equation

We have differential form of Gauss theorem, $\nabla \cdot \vec{E} = \rho / \epsilon_0$ — (1)

Again we know, $\vec{E} = -\nabla V$ — (ii)

Now, from eq^s (1) & (ii), $\nabla \cdot (-\nabla V) = \rho / \epsilon_0$

$$\Rightarrow \nabla^2 V = -\rho / \epsilon_0 \text{ — Poisson's eqⁿ}$$

Now, for field free region $\rho = 0$

$$\therefore \nabla^2 V = 0 \text{ — This is Laplace's equation}$$

Now, $|\vec{r}-\vec{d}|^2 = |\vec{r}-\vec{d}| \cdot |\vec{r}-\vec{d}| \stackrel{(8)}{=} \vec{r} \cdot \vec{r} - \vec{r} \cdot \vec{d} - \vec{d} \cdot \vec{r} + \vec{d} \cdot \vec{d}$
 $= r^2 + d^2 - 2rd \cos \theta = r^2 \left[1 + \frac{d^2}{r^2} - \frac{2d \cos \theta}{r} \right]$
 $= r^2 \left[1 - \frac{2d \cos \theta}{r} \right]$

$|\vec{r}-\vec{d}| = r \left[1 - \frac{2d \cos \theta}{r} \right]^{1/2}$ $\frac{d}{r} \rightarrow \text{small}$
 $\frac{d^2}{r^2} \rightarrow 0$

$\frac{1}{|\vec{r}-\vec{d}|} \approx \frac{1}{r} \left[1 + \frac{d \cos \theta}{r} \right] = \frac{1}{r} \left[1 + \frac{d \cdot \vec{r} \cdot \vec{d}}{r^2} \right]$

Again from eq (1) $V(r, \theta) = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{r} + \frac{d \cdot \vec{r}}{r^3} - \frac{1}{r} \right]$

$V(r, \theta) = \frac{q}{4\pi\epsilon_0} \frac{d \cdot \vec{r}}{r^3}$

Again, $\vec{\nabla} \left(\frac{1}{r} \right) = -\frac{\vec{r}}{r^3}$ $V(r, \theta) = -\frac{q}{4\pi\epsilon_0} \vec{\nabla} \left(\frac{1}{r} \right) \cdot \vec{d}$ $\vec{M} = q\vec{d}$

Now, $E(r, \theta) = -\vec{\nabla} V(r, \theta) = -\vec{\nabla} \left[\frac{q}{4\pi\epsilon_0} \frac{d \cdot \vec{r}}{r^3} \right]$

$= -\frac{q}{4\pi\epsilon_0} \left[(\vec{d} \cdot \vec{r}) \vec{\nabla} \left(\frac{1}{r^3} \right) + \frac{1}{r^3} \vec{\nabla} (\vec{d} \cdot \vec{r}) \right]$

we have, $\vec{\nabla} \left(\frac{1}{r^3} \right) = -\frac{3\vec{r}}{r^5}$
 $\vec{\nabla} (\vec{d} \cdot \vec{r}) = \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) (d_x \hat{i} + d_y \hat{j} + d_z \hat{k}) = \vec{d}$

$E(r, \theta) = -\frac{q}{4\pi\epsilon_0} \left[(\vec{d} \cdot \vec{r}) \left(-\frac{3\vec{r}}{r^5} \right) + \frac{1}{r^3} \vec{d} \right]$

$E(r, \theta) = \frac{q}{4\pi\epsilon_0} \left[\frac{3(\vec{d} \cdot \vec{r}) \vec{r}}{r^5} - \frac{\vec{d}}{r^3} \right]$

We have $E(r, \theta) = -\vec{\nabla} V$ (9)

Two dimensional polar coordinate

$$E(r, \theta) = - \left[\frac{\partial}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial}{\partial \theta} \hat{\theta} \right] \left(\frac{M \omega \sin \theta}{4\pi \epsilon_0 r} \right)$$

$$= - \frac{M}{4\pi \epsilon_0} \left[- \frac{2\omega \sin \theta}{r^3} \hat{r} - \frac{\cos \theta}{r^3} \hat{\theta} \right]$$

$$E(r, \theta) = + \frac{M}{4\pi \epsilon_0 r^3} \left[2\omega \sin \theta \hat{r} + \cos \theta \hat{\theta} \right]$$

$$E_r \hat{r} + E_\theta \hat{\theta} = \frac{M}{4\pi \epsilon_0 r^3} \left[2\omega \sin \theta \hat{r} + \cos \theta \hat{\theta} \right]$$

$$= \frac{2M \omega \sin \theta}{4\pi \epsilon_0 r^3} \hat{r} + \frac{M \cos \theta}{4\pi \epsilon_0 r^3} \hat{\theta}$$

∴ Radial Component of electric field $E_r = \frac{2M \omega \sin \theta}{4\pi \epsilon_0 r^3}$

Angular $E_\theta = \frac{M \cos \theta}{4\pi \epsilon_0 r^3}$

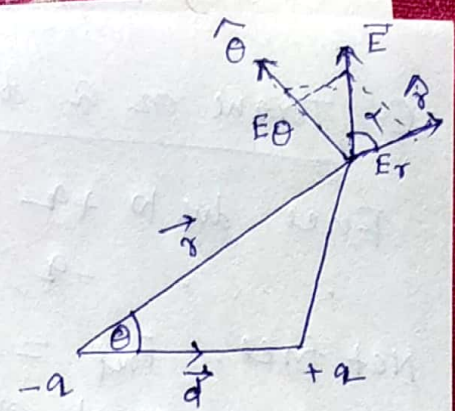
$$|\vec{E}| = \sqrt{(E_r)^2 + (E_\theta)^2} = \sqrt{\left(\frac{2M \omega \sin \theta}{4\pi \epsilon_0 r^3} \right)^2 + \left(\frac{M \cos \theta}{4\pi \epsilon_0 r^3} \right)^2}$$

$$= \frac{M}{4\pi \epsilon_0 r^3} \sqrt{4\omega^2 \sin^2 \theta + \cos^2 \theta}$$

$$|\vec{E}| = \frac{M}{4\pi \epsilon_0 r^3} \sqrt{3\omega^2 \sin^2 \theta + 1}$$

$$\tan \alpha = \frac{E_\theta}{E_r} = \frac{M \cos \theta}{4\pi \epsilon_0 r^3} \times \frac{4\pi \epsilon_0 r^3}{2M \omega \sin \theta} = \frac{1}{2} \cot \theta$$

$$\alpha = \tan^{-1} \left(\frac{1}{2} \cot \theta \right)$$



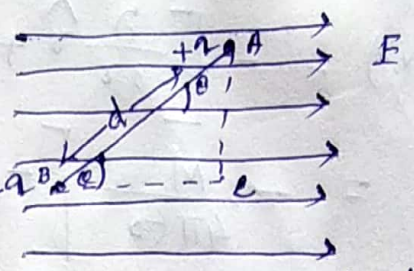
$$V(r, \theta) = \frac{M \omega \sin \theta}{4\pi \epsilon_0 r^2}$$

Torque on a dipole in uniform electric field

Force due to +q $F_1 = qE$ — (1)

Force due to -q $F_2 = -qE$ — (2)

Net force, $F_{net} = 0$



But two equal and opposite force constitute a couple. Due to this couple tends rotate the electric dipole clockwise and try to align to the electric field

Now torque acting on the dipole = magnitude of the either force x Perpⁿ distance betⁿ two forces.

$$= qE \times AC$$

from ΔABC , $\sin\theta = \frac{AC}{AB} \Rightarrow AC = d \sin\theta$

$$= qE \times d \sin\theta = ME \sin\theta$$

$\vec{\tau} = \vec{M} \times \vec{E}$

① uniqueness theorem :-

two solutions of Laplace's equation observing same boundary condition differ by a const.

i.e. $\phi_1 - \phi_2 = C$

let ϕ_1 & ϕ_2 two solⁿ of Laplace's eqⁿ observing boundary condition

(i) Dirichlet condition, i.e. $\phi_1 = \phi_2$

(ii) Neumann " i.e. $\frac{\partial \phi_1}{\partial n} = \frac{\partial \phi_2}{\partial n}$

let, $\phi = \phi_1 - \phi_2$
 $\nabla^2 \phi = \nabla^2 \phi_1 - \nabla^2 \phi_2 \Rightarrow$

i.e. $\nabla^2 \phi = 0$

$\nabla^2 \phi_1 - \nabla^2 \phi_2 \Rightarrow$

$\phi_1 = \phi_2$